

# 修正的CH方程的非零渐近值的光滑孤立波的直接求法

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## 摘要

本文关注mCH方程的非零渐近值的光滑孤立波, 并通过平面动力系统的分析方法, 直接给出了孤立波解的显示表达式。

## 关键词

修正的Camassa-Holm方程, 孤立波解, 平面动力系统方法

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# Direct Method for Smooth Solitary Waves of Non-Zero Asymptotic Values of the Modified CH Equation

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## Abstract

This paper focuses on the non-zero asymptotic value of the smooth solitary wave of the mCH equation and directly gives the explicit expression of the solitary wave solution through the analysis of the planar dynamical system.

## Keywords

Modified Camsa-Holm Equation, Solitary Wave Solution, Planar Dynamical System Method

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## 1. 背景

著名的 Camassa-Holm (CH) 方程的形式是

$$u_t - u_{xxt} + 2ku_x + 3uu_x = 2u_xu_{xx} + uu_{xxx}. \quad (1)$$

方程(1)在 1981 年被 Fuchssteiner 和 Fokas [1] 引入作为一种新的可积系统。后来, Camassa 和 Holm [2] [3] 发现了方程式可以作为浅水波模型, 并且当  $k=0$  时, 它具有非光滑孤立波, 表达式为

$$u(x, t) = ce^{-|x-ct|},$$

其中  $c$  是波速。进一步的, Constantin and Strauss [4] 证明了上述解是轨道稳定的。

当  $k \neq 0$  时, 刘[5]等人证明了方程(1)具有非零渐近值的非光滑孤立波, 表达式为

$$u(x, t) = (k+c)e^{-|x-ct|} - k.$$

欧阳等人证明了该解是轨道稳定的。许多其他的作者都研究了方程(1)的很多其他性质[6]-[12]。

此外, 方程(1)的其他形式都被广泛的研究, 例如, Li and Olver [13] 考虑了方程

$$u_t - u_{xxt} + 2ku_x + auu_x = 2u_xu_{xx} + uu_{xxx} \quad (2)$$

的全局适定性, 其中  $a$  是常数。基于方程(2), 刘[14]等人推广了该方程, 修正的 CH 方程的具体形式为

$$u_t - u_{xxt} + 2ku_x + au^2u_x = 2u_xu_{xx} + uu_{xxx}. \quad (3)$$

部分学者研究了该方程的很多性质[15] [16] [17] [18]。其中当  $k=0$  时, Wazwaz [19] 研究了方程(3)的钟型孤立波解, 表达式为

$$u(x, t) = -2 \operatorname{sech}^2 \frac{1}{2}(x-ct).$$

在这篇文章中, 主要关注非零渐近值的光滑孤立波解, 并采用平面动力系统的分析方法, 直接求出其孤立波解的精确表达式。这个方法同样适用于求解其他解, 例如周期波解, 扭波解, 尖孤立波解等等。

## 2. 平面动力系统

首先, 令  $u(x, t) = \varphi(\xi)$ , 且  $\xi = x - ct$  代入方程(3), 得到常微分方程

$$-c\varphi' + c\varphi''' + 2k\varphi' + a\varphi^2\varphi' = 2\varphi'\varphi'' + \varphi\varphi'''. \quad (4)$$

积分方程(4), 得到

$$\varphi''(\varphi - c) = \frac{a}{3}\varphi^3 + (2k - c)\varphi - \frac{1}{2}(\varphi')^2 + g,$$

其中  $g$  是一个积分常数。

让  $y = \varphi'$ , 我们将得到下面的平面系统

$$\begin{cases} \varphi' = y \\ y' = \frac{1}{\varphi - c} \left( \frac{a}{3}\varphi^3 + (2k - c)\varphi - \frac{1}{2}y^2 + g \right). \end{cases}$$

两边同时乘以  $\varphi - c$  得到

$$\begin{cases} (\varphi - c)\varphi' = (\varphi - c)y \\ (\varphi - c)y' = \left( \frac{a}{3}\varphi^3 + (2k - c)\varphi - \frac{1}{2}y^2 + g \right) \end{cases}$$

做变换  $d\xi = (\varphi - c)d\tau$ , 我们就可以得到如下的哈密尔顿系统

$$\begin{cases} \frac{d\varphi}{d\tau} = (\varphi - c)y \\ \frac{dy}{d\tau} = \frac{a}{3}\varphi^3 + (2k - c)\varphi - \frac{1}{2}y^2 + g \end{cases}$$

上述系统具有首次积分

$$H(\varphi, y) = h,$$

其中

$$H(\varphi, y) = y^2(\varphi - c) - \frac{a}{6}\varphi^4 - (2k - c)\varphi^2 - 2gy\varphi.$$

### 3. 孤立波解的直接求法

令

$$f(\varphi) = \frac{a}{6}(\varphi - c)(\varphi - \alpha)^2(\varphi - \beta),$$

其中,  $\alpha$  和  $\beta$  是  $f(\varphi) = 0$  的两个根。进一步的, 从首次积分中我们可以得到

$$\alpha = \frac{-ac + \Delta}{3a},$$

$$\beta = \frac{-ac - 2\Delta}{3a},$$

其中  $\Delta = \sqrt{-2a^2c^2 - 18a(2k - c)}$ , 因此, 首次积分可以写成

$$y^2 = \frac{a}{6}(\varphi - \alpha)^2(\varphi - \beta),$$

或者

$$y = \pm \sqrt{\frac{a}{6}} \sqrt{\varphi - \beta} (\varphi - \alpha).$$

将上述  $y$  的表达式代入到  $\frac{d\varphi}{d\xi} = y$  中, 并且积分一次得到

$$\int_{\beta}^{\varphi} \frac{ds}{\sqrt{s - \beta}(s - \alpha)} = \sqrt{\frac{a}{6}} \int_0^{\xi} d\xi,$$

完成这个积分我们就得到

$$\varphi(\xi) = \alpha + (\beta - \alpha) \operatorname{sech}^2 \left( \frac{\eta}{2} \right),$$

其中

$$\eta = \sqrt{\frac{a(\alpha - \beta)}{6}} \xi,$$

注意到  $u(x, t) = \varphi(\xi)$ ，所以我们得到孤立波解为

$$u(x, t) = \alpha + (\beta - \alpha) \operatorname{sech}^2\left(\frac{\eta}{2}\right).$$

这个解具有非零渐近值  $\alpha$ 。这个解可以通过数学软件 Mathematica 验证其正确性。其次若参数取特定的值的时候，这个解与文献[19]的解是一致的，进一步验证了该解的正确性。

## 4. 结论

该文章所使用的动力系统的分支方法不仅适用于孤立波解，也同样适用于行波解，周期波解以及不光滑的孤立波解。分析的过程类似，需要克服的困难是选定特定的初值和积分区域，从而得到的解就不同。与齐次平衡法，待定系数法，扰动方法等比较，我们的方法可以求出具有 Hamiltonian 函数的所有有界行波解。这是动力系统分支方法的独特之处。接下来，我们还将应用于其他更多的具有力学，物理学等背景的偏微分方程中。

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