

New Exact Analytic Solutions to a Coupled KdV Equations with Variable Coefficients*

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Abstract: Based on the homogeneous balance principle and general variable separation approach, with the aid of two generalized Riccati equations and Mathematica software, we first find some exact analytic solutions to a coupled KdV equations with variable coefficients, including several kinds of soliton-like solutions, periodical-like solutions and solitary wave solutions with variable speed. Some of them are found for the first time.

Keywords: Coupled KdV Equations with Variable Coefficients; Exact Analytic Solutions; Riccati Equations; Solitary Wave Solutions with Variable Speed

一类变系数组合 KdV 方程新的精确解析解*

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摘要: 在齐次平衡法和分离变量的基础上, 通过两个推广的 Riccati 方程, 借助 Mathematica 软件, 求出了一类变系数组合 KdV 方程的一些精确解析解, 包括各种类孤立子解、类周期解和变速孤立波解, 部分解为首次发现。

关键词: 变系数组合 KdV 方程; 精确解析解; Riccati 方程; 变速孤立波

1. 引言

非线性发展方程的求解是广大物理学、力学、地球科学、生命科学、应用数学、和工程技术工作者研究的一个重要课题, 多年来众多科学家为此做了大量的工作。常系数非线性方程只能近似地反映实际物质运动变化规律, 而变系数非线性方程却能更加准确地描述物质的属性, 因此研究变系数非线性方程的精确解显得十分重要, 近年来, 人们已经发现了一些有效的求解方法, 如: 反散射法^[1]、Backlund 变换法^[2]、Darboux 变换法^[3]、Jacobi 椭圆函数法^[4]等等^[5,6]。本文运用两个推广形式的 Riccati 方程组^[7-8]求解一类变系数组合 KdV 方程

$$\begin{cases} u_t + \alpha(t)uu_x + \beta(t)vv_x + \gamma(t)u_{xxx} = 0 \\ v_t + \delta(t)uv_x + \gamma(t)v_{xxx} = 0 \end{cases} \quad (1)$$

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方程(1)是 Hirota-Satsuma 组合 KdV 方程^[9]的推广, 当 $v = v(x, t) = 0$ 时, (1)转化为变系数 KdV 方程^[10-11], 文献[12]利用 F 展开法研究了方程(1)的 Jacobi 椭圆函数解。本文将利用两个推广的 Riccati 方程研究(1)的精确解, 得到了许多新的结果。

2. 推广的 Riccati 方程法

对非线性发展方程:

$$P(u, u_t, u_x, uu_x, u_{xt}, u_{xx}, u_{xxx}, \dots) = 0 \quad (2)$$

我们寻求如下形式的解:

$$u(x, t) = \sum_{i=0}^n a_i f^i(\xi) + \sum_{j=1}^n b_j f^{j-1}(\xi) g(\xi) \quad (3)$$

其中 $a_0 = a_0(t), a_i = a_i(t), b_j = b_j(t), (i, j = 1, 2, \dots, n)$, $\xi = \xi(x, t)$ 均是关于相应变元的任意函数, n 是待定常数, 它可以通过平衡最高阶导数项和非线性项确定,

而 $f(\xi)$ 、 $g(\xi)$ 满足如下投影 Riccati 方程组:

$$(1) \begin{cases} f'(\xi) = -qf(\xi)g(\xi), \\ g'(\xi) = q[1 - g^2(\xi) - rf(\xi)], \\ g^2(\xi) = 1 - 2rf(\xi) + (r^2 + \varepsilon)f^2(\xi) \end{cases} \quad (4)$$

这里 ' 表示 $\frac{d}{d\xi}$, r, q 为任何实数, $\varepsilon = \pm 1$, 后面雷同,

方程组(4)有下列解:

$$\begin{cases} f_1(\xi) = \frac{a}{b \cosh(q\xi) + c \sinh(q\xi) + ar}, \\ g_1(\xi) = \frac{b \sinh(q\xi) + c \cosh(q\xi)}{b \cosh(q\xi) + c \sinh(q\xi) + ar} \end{cases} \quad (5)$$

其中 a, b, c 满足条件: 当 $\varepsilon = 1$ 时 $c^2 = a^2 + b^2$, 当 $\varepsilon = -1$ 时 $b^2 = a^2 + c^2$ 。

$$(2) \begin{cases} f'(\xi) = qf(\xi)g(\xi), \\ g'(\xi) = q[1 + g^2(\xi) - rf(\xi)], \\ g^2(\xi) = -1 + 2rf(\xi) + (1 - r^2)f^2(\xi) \end{cases} \quad (6)$$

方程组(6)有下列解:

$$\begin{cases} f_2(\xi) = \frac{a}{b \cos(q\xi) + c \sin(q\xi) + ar}, \\ g_2(\xi) = \frac{b \sin(q\xi) - c \cos(q\xi)}{b \cos(q\xi) + c \sin(q\xi) + ar} \end{cases} \quad (7)$$

$$(1) \alpha(t) = c_1\delta(t), \beta(t) = c_2\delta(t), \gamma(t) = c_3\delta(t)$$

$$\begin{aligned} a_1 = a_3 = a_4 = 0, b_1 = b_3 = b_4 = 0, r = 0, \varepsilon = \pm 1, \omega = -(a_0k + 4k^3q^2c_3) \int \delta(t) dt \\ a_2 = -12\varepsilon k^2 q^2 c_3, b_2 = \mp 12\varepsilon k^2 q^2 c_3 \sqrt{\frac{1-c_1}{c_2}}, b_0 = \pm \frac{a_0 + a_0c_1 - 16k^2 q^2 c_3}{\sqrt{c_2 - c_1c_2}} \end{aligned} \quad (10)$$

其中 $c_2 - c_1c_2 \neq 0$, c_1, c_2, c_3, k, q, a_0 均为任意常数, 后面雷同。

$$(2) \gamma(t) = c_3\delta(t), \alpha(t) = c_1\delta(t) + c_2\beta(t)$$

$$\begin{aligned} a_2 = a_3 = a_4 = 0, b_2 = b_3 = b_4 = 0, r = \pm\sqrt{-\varepsilon}, \varepsilon = \pm 1, \omega = -(a_0k + k^3q^2c_3) \int \delta(t) dt \\ a_1 = \pm 6k^2 q^2 c_3 \sqrt{-\varepsilon}, b_1 = \pm 6k^2 q^2 c_3 \sqrt{\frac{\delta(t) - c_1\delta(t) - c_2\beta(t)}{-\varepsilon\beta(t)}}, b_0 = \pm a_0 \sqrt{\frac{\delta(t) - c_1\delta(t) - c_2\beta(t)}{\beta(t)}} \end{aligned} \quad (11)$$

由(5)(8)(9)(10)和(5)(8)(9)(11)我们可以得到下列解

$$\begin{aligned} u_1(\xi_1) &= a_0 + \frac{-12\varepsilon k^2 q^2 c_3 a^2}{[b \cosh(q\xi_1) + c \sinh(q\xi_1)]^2}, \\ v_1(\xi_1) &= \pm \frac{a_0 + a_0c_1 - 16k^2 q^2 c_3}{\sqrt{c_2 - c_1c_2}} \mp \frac{12\varepsilon k^2 q^2 a^2 c_3}{[b \cosh(q\xi_1) + c \sinh(q\xi_1)]^2} \sqrt{\frac{1-c_1}{c_2}} \\ \xi_1 &= \xi_1(x, t) = kx - (a_0k + 4k^3q^2c_3) \int \delta(t) dt + \xi_0 \\ u_2(\xi) &= a_0 \pm \frac{6k^2 q^2 c_3 \sqrt{-\varepsilon}}{b \cosh(q\xi_2) + c \sinh(q\xi_2) \pm a\sqrt{-\varepsilon}}, \end{aligned}$$

其中 a, b, c 满足条件: $a^2 = b^2 + c^2$ 。

将(3), (4)和(3), (6)分别代入(1)并令 $f^i(\xi)g^j(\xi)$ 系数为零 ($i = 1, 2, \dots; j = 0, 1, \dots$), 可得一关于所有待定系数的非线性代数方程组(NAEs), 借助 Mathematica 软件求解该 NAEs 便可由(5)(7)得(1)的精确解。

3. 方程(1)的精确解

由齐次平衡原则可设

$$u(\xi) = a_0 + a_1f(\xi) + a_2f^2(\xi) + a_3f(\xi)g(\xi) + a_4g(\xi) \quad (8)$$

$$v(\xi) = b_0 + b_1f(\xi) + b_2f^2(\xi) + b_3f(\xi)g(\xi) + b_4g(\xi) \quad (9)$$

其中 $\xi = \xi(x, t) = kx + \omega(t) + \xi_0$, $\omega(t)$, k, a_i, b_i ($i = 1, 2, 3, 4$) 待定, ξ_0 为任意常数。

情形 1

将(4)(8)(9)代入(1)并令 $f^i(\xi)g^j(\xi)$ ($i = 1, 2, \dots; j = 0, 1$) 系数为零, 得一关于 $a_0, a_1, a_2, a_3, a_4, b_0, b_1, b_2, b_3, b_4, k, \omega(t)$ 的代数方程组, 借助 mathematica 和吴消元法可得解:

$$v_2(\xi_2) = \pm a_0 \sqrt{\frac{\delta(t) - c_1 \delta(t) - c_2 \beta(t)}{\beta(t)}} \pm \frac{6k^2 q^2 c_3 a \sqrt{\frac{\delta(t) - c_1 \delta(t) - c_2 \beta(t)}{-\varepsilon \beta(t)}}}{b \cosh(q\xi_2) + c \sinh(q\xi_2) \pm a\sqrt{-\varepsilon}}$$

其中 $\xi_2 = \xi_2(x, t) = kx - (a_0 k + k^3 q^2 c_3) \int \delta(t) dt + \xi_0$

情形 2

同理将(6)、(8)(9)代入(1)并令

$a_0, a_1, a_2, a_3, a_4, b_0, b_1, b_2, b_3, b_4, k, \omega(t)$ 的代数方程组,

$f^i(\xi) g^j(\xi)$ ($i=1, 2, \dots; j=0, 1$) 系数为零, 得一关于

借助 mathematica 和吴消元法可得解:

(3) $\alpha(t) = c_1 \delta(t), \beta(t) = c_2 \delta(t), \gamma(t) = c_3 \delta(t)$

$$\begin{aligned} a_1 = a_3 = a_4 = 0, b_1 = b_3 = b_4 = 0, r = 0, \omega = (-a_0 k + 4k^3 q^2 c_3) \int \delta(t) dt \\ a_2 = 12k^2 q^2 c_3, b_2 = \mp 12k^2 q^2 c_3 \sqrt{\frac{-1-c_1}{c_2}}, b_0 = \pm \frac{-a_0 + a_0 c_1}{\sqrt{-c_2 - c_1 c_2}} \end{aligned} \tag{12}$$

(4) $\gamma(t) = c_3 \delta(t), \beta(t) = c_1 \delta(t) + c_2 \alpha(t)$

$$\begin{aligned} a_2 = a_3 = a_4 = 0, b_2 = b_3 = b_4 = 0, r = \pm 1, \omega = (-a_0 k + k^3 q^2 c_3) \int \delta(t) dt \\ a_1 = -6k^2 q^2 r c_3, b_1 = \mp 6k^2 q^2 r c_3 \sqrt{\frac{-\alpha(t) + \delta(t)}{c_1 \delta(t) + c_2 \alpha(t)}}, b_0 = \pm a_0 \sqrt{\frac{-\alpha(t) + \delta(t)}{c_1 \delta(t) + c_2 \alpha(t)}} \end{aligned} \tag{13}$$

由(7)(8)(9)(12)和(7)(8)(9)(13)我们可以得到下列解

$$\begin{aligned} u_3(\xi_3) &= a_0 + \frac{12k^2 q^2 c_3 a^2}{[b \cos(q\xi_3) + c \sin(q\xi_3)]^2}, \\ v_3(\xi_3) &= \pm \frac{-a_0 + a_0 c_1}{\sqrt{-c_2 - c_1 c_2}} \mp \frac{12k^2 q^2 a^2 c_3 \sqrt{\frac{-1-c_1}{c_2}}}{[b \cos(q\xi_3) + c \sin(q\xi_3)]^2} \\ \xi_3 &= \xi_3(x, t) = kx + (-a_0 k + 4k^3 q^2 c_3) \int \delta(t) dt + \xi_0 \\ u_4(\xi_4) &= a_0 - \frac{6k^2 q^2 r c_3 a}{b \cos(q\xi_4) + c \sin(q\xi_4) + ar}, \\ v_4(\xi_4) &= \pm a_0 \sqrt{\frac{-\alpha(t) + \delta(t)}{c_1 \delta(t) + c_2 \alpha(t)}} \mp \frac{6k^2 q^2 r a c_3 \sqrt{\frac{-\alpha(t) + \delta(t)}{c_1 \delta(t) + c_2 \alpha(t)}}}{b \cos(q\xi_4) + c \sin(q\xi_4) + ar} \\ \xi_4 &= \xi_4(x, t) = kx + (-a_0 k + k^3 q^2 c_3) \int \delta(t) dt + \xi_0 \end{aligned}$$

若在 u_1, v_1 和 u_3, v_3 中取 $c = 0, b = a$ 或 $b = 0, c = a$, 我们可以得到类钟型孤子解、类奇异孤立波解和类周期解:

$$\begin{aligned} u_{11}(\xi_{11}) &= a_0 - 12\varepsilon k^2 q^2 c_3 \operatorname{sech}^2(q\xi_{11}), \\ v_{11}(\xi_{11}) &= \pm \frac{a_0 + a_0 c_1 - 16k^2 q^2 c_3}{\sqrt{c_2 - c_1 c_2}} \mp 12\varepsilon k^2 q^2 c_3 \sqrt{\frac{1-c_1}{c_2}} \operatorname{sech}^2(q\xi_{11}) \\ u_{12}(\xi_{12}) &= a_0 - 12\varepsilon k^2 q^2 c_3 \operatorname{csch}^2(q\xi_{12}), \\ v_{12}(\xi_{12}) &= \pm \frac{a_0 + a_0 c_1 - 16k^2 q^2 c_3}{\sqrt{c_2 - c_1 c_2}} \mp 12\varepsilon k^2 q^2 c_3 \sqrt{\frac{1-c_1}{c_2}} \operatorname{csch}^2(q\xi_{12}) \end{aligned}$$

$$u_{31}(\xi_{31}) = a_0 + 12k^2 q^2 c_3 \sec^2(q\xi_{31}),$$

$$v_{31}(\xi_{31}) = \pm \frac{-a_0 + a_0 c_1}{\sqrt{-c_2 - c_1 c_2}} \mp 12k^2 q^2 c_3 \sqrt{\frac{-1 - c_1}{c_2}} \sec^2(q\xi_{31})$$

$$u_{32}(\xi_{32}) = a_0 + 12k^2 q^2 c_3 \csc^2(q\xi_{32}),$$

$$v_{32}(\xi_{32}) = \pm \frac{-a_0 + a_0 c_1}{\sqrt{-c_2 - c_1 c_2}} \mp 12k^2 q^2 c_3 \sqrt{\frac{-1 - c_1}{c_2}} \csc^2(q\xi_{32})$$

其中 $\xi_{11} = \xi_{12} = kx - (a_0 k + 4k^3 q^2 c_3) \int \delta(t) dt + \xi_0$,

$\xi_{31} = \xi_{32} = kx - (a_0 k - 4k^3 q^2 c_3) \int \delta(t) dt + \xi_0$

这里的波速是 t 的函数, 随 t 的改变而改变。

注记: 解 $u_{11}, v_{11}, u_{31}, v_{31}$ 包含文献[12]中的解(3.25)(3.26)(3.31)(3.32), 若选择 a, b, c 不同的取值, 我们可以得到方程(1)其他新的精确解, 据我们所知, 这里的四组解是新的。

4. 结论

本文运用两个推广的 Riccati 方程组, 求解一类变系数组合 KdV 方程, 得到了该方程的一些精确解析解, 包括类孤立波解, 类周期解和变速孤立波解。实践证明这种方法可以适用于许多其它非线性方程, 如何将该方法推广到具高次非线性耦合及离散方程, 还值得进一步研究。

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